

Computing Interior Eigenvalues of a Generalized Complex-Symmetric Pencil arising from the Modeling of Resonant MEMS Systems

Tsuyoshi Koyama

Sanjay Govindjee

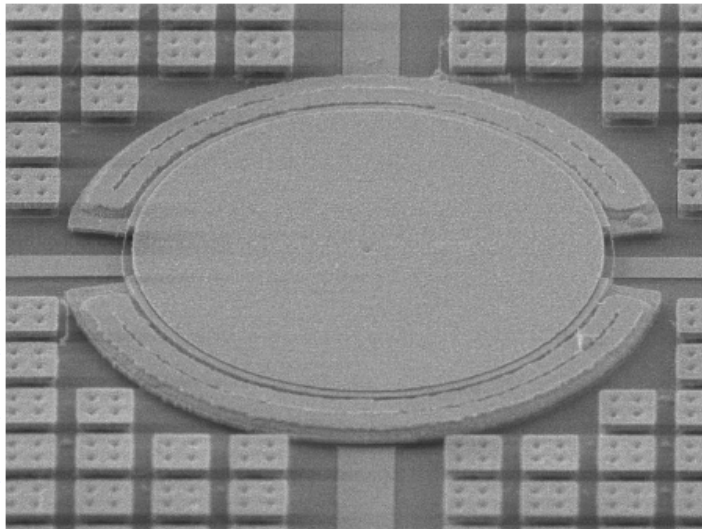
Civil Engineering, UC Berkeley

Center for Mechanics, IMES, ETHZ

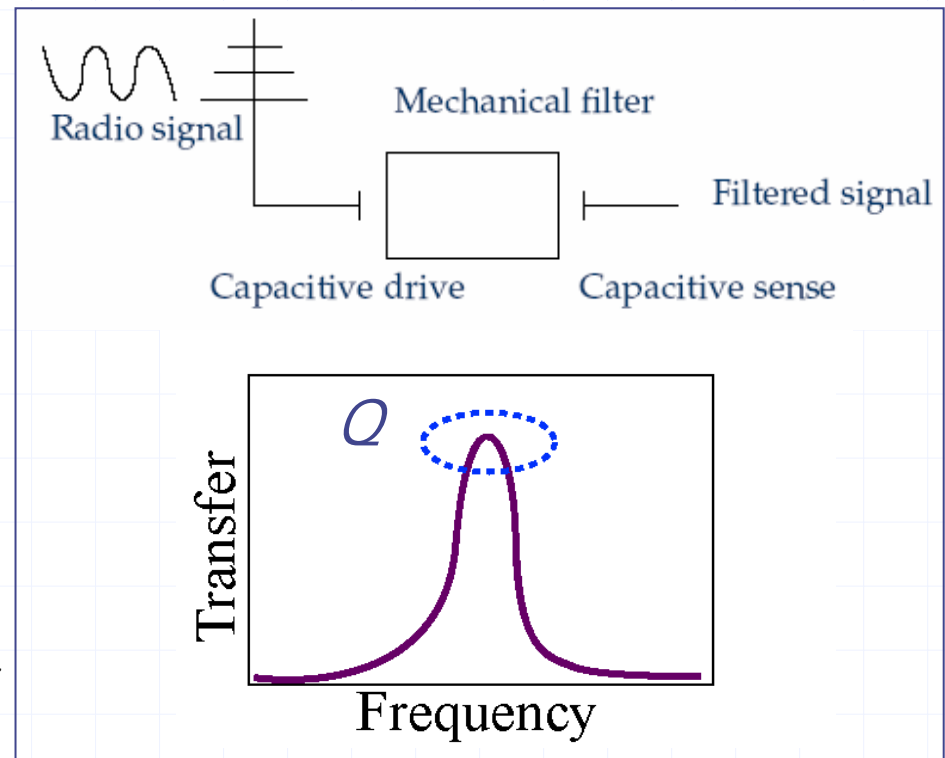


High-frequency MEMS resonators (MHz-GHz)

- ◆ Applications as small size, low energy consuming frequency references, filters, and sensors



SEM of 41.5 μm radius poly-SiGe disk resonator



Resonator simulation

◆ Design requires knowledge of

- Frequency

- **Quality factor (Q)**

Existing
Software
and Methods

?

$$Q = \frac{\text{Maximum Stored Energy}}{\text{Energy Loss per radian}} \approx \frac{1}{\text{Damping}}$$



Tools and methods for evaluating damping in resonant MEMS

Mathematical problem

- ◆ Equation of motion discretized with FEM under harmonic assumption

Quadratic eigenvalue problem.

$$\left(-\omega^2 \mathbf{M} + \mathbf{K}\right) \mathbf{x} = \mathbf{0}$$

Complex eigenfrequency

$$\omega = \omega_R + i\omega_I$$

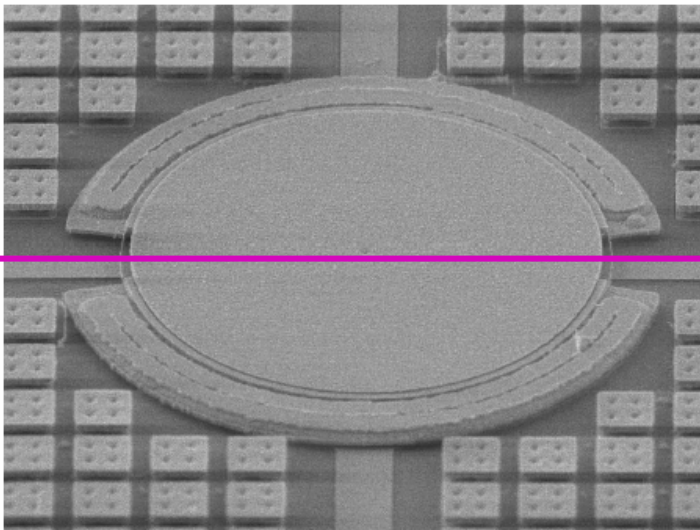
$$Q = \frac{|\omega|}{2\omega_I}$$

Overview

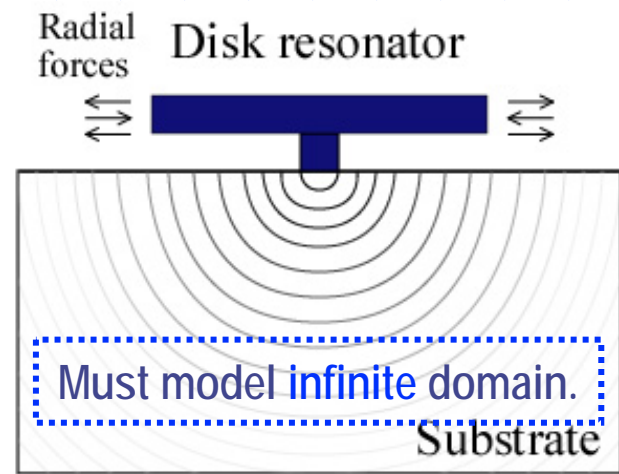
- ◆ Characteristics of the eigenvalue problem
 - Complex symmetry from Perfectly Matched Layers
- ◆ Solution method
 - Projection methods
 - Jacobi-Davidson QZ method (JDQZ)
 - Geometric multigrid preconditioned GMRES for the solution of the correction equation
- ◆ Numerical example of a disk resonator
- ◆ Conclusions

Damping mechanism: Anchor loss

- ◆ Mechanism: Energy loss from radiating waves escaping into the substrate.



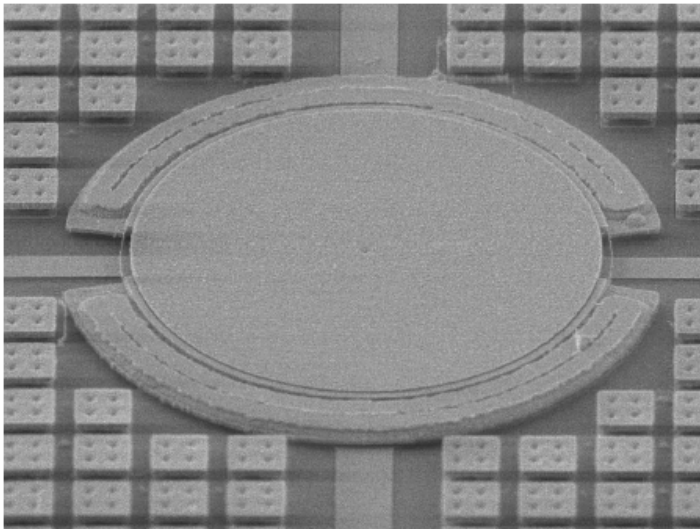
SEM of 41.5 μm radius poly-SiGe disk resonator



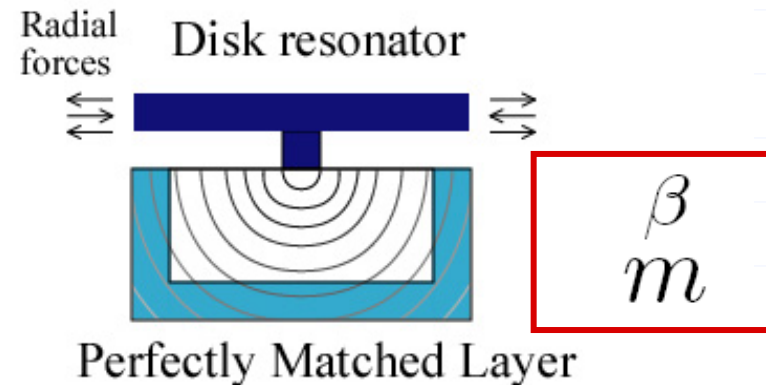
Section of disk resonator

Perfectly Matched Layers (PML)

- ◆ Mechanism: Energy loss from radiating waves escaping into the substrate.



SEM of 41.5 μm radius poly-SiGe disk resonator



Outgoing waves are absorbed with **zero impedance** mismatch at PML boundaries.

Difficulties

◆ PML implementation:

Generalized eigenvalue problem

$$\mathbf{K}\mathbf{x} = \omega^2 \mathbf{M}\mathbf{x}$$

1. Sparse matrices are large (millions and larger) for accurate results
Requires large scale sparse-eigenvalue methods.
2. Matrices are complex valued symmetric, Non-Hermitian
Complex-valued eigenvalues,
3. Desired modes are not at the exterior of the spectrum but interior
Adds difficulty in the linear solves. Little knowledge about optimal iterative linear system solution methods. (for complex symmetric systems).

Proposed solution

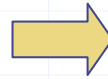
1. Sparse matrices are large (millions and larger) for accurate results

Projection methods.

2. Matrices are complex valued, Non-Hermitian

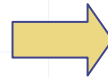
3. Desired modes are not at the exterior of the spectrum but interior

Elastodynamic equation
(SPD pencil without PML)



Geometric multigrid
preconditioned GMRES
linear solver

Iterative solvers become
expensive for higher accuracy



Jacobi-Davidson QZ
eigenvalue method

Projection methods

- ◆ Find eigen pair $(\omega^2, \mathbf{x} \in \mathcal{V})$ such that

$$(\mathbf{K} - \omega^2 \mathbf{M}) \mathbf{x} \perp \mathcal{W}$$

- ◆ Galerkin projection (Ritz values, vectors)
- ◆ Petrov-Galerkin projection

$$\begin{aligned} \mathcal{W} &= \mathcal{V} \\ \mathcal{W} &\neq \mathcal{V} \end{aligned}$$

Select harmonic-Ritz projection for better approximation of interior eigenvalues.

$$\mathcal{W} = (\mathbf{K} - \omega_0^2 \mathbf{M}) \mathcal{V}$$

Jacobi-Davidson QZ (JDQZ)

◆ Harmonic JD method for (\mathbf{K}, \mathbf{M}) target ω_0^2

Given subspace \mathcal{V} , $\mathcal{W} = (\mathbf{K} - \omega_0^2 \mathbf{M}) \mathcal{V}$

1. Find approximate eigenpair (ω^2, \mathbf{u})

$$\mathbf{u} \in \mathcal{V} \quad (\mathbf{K} - \omega^2 \mathbf{M}) \mathbf{u} \perp \mathcal{W}$$

2. Solve **correction equation** for \mathbf{t}

3. Expand subspace

$$\begin{aligned}\mathcal{V}_{\text{new}} &= \mathcal{V} \oplus \text{span}\{\mathbf{t}\} \\ \mathcal{W}_{\text{new}} &= \mathcal{W} \oplus \text{span}\{(\mathbf{K} - \omega_0^2 \mathbf{M}) \mathbf{t}\}\end{aligned}$$

(Fokkema, Sleijpen, and van der Horst, 1998)

Correction equation

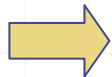
- ◆ Find $\mathbf{t} \perp \mathbf{u}$ such that $(\mathbf{u} + \mathbf{t})$ is a better approximation to the desired eigenvector

$$(\mathbf{I} - \mathbf{p}\mathbf{p}^*) (\mathbf{K} - \omega^2 \mathbf{M}) (\mathbf{I} - \mathbf{u}\mathbf{u}^*) \mathbf{t} = -\mathbf{r}$$

$$\begin{aligned} \mathbf{r} &= (\mathbf{K} - \omega^2 \mathbf{M}) \mathbf{u} \\ \mathbf{p} &= (\mathbf{K} - \omega_0^2 \mathbf{M}) \mathbf{u} \in \mathcal{W} \end{aligned}$$

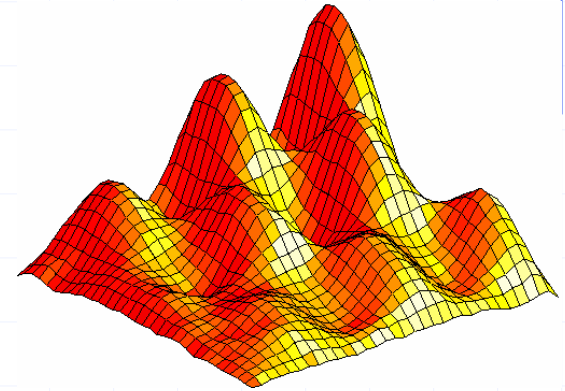
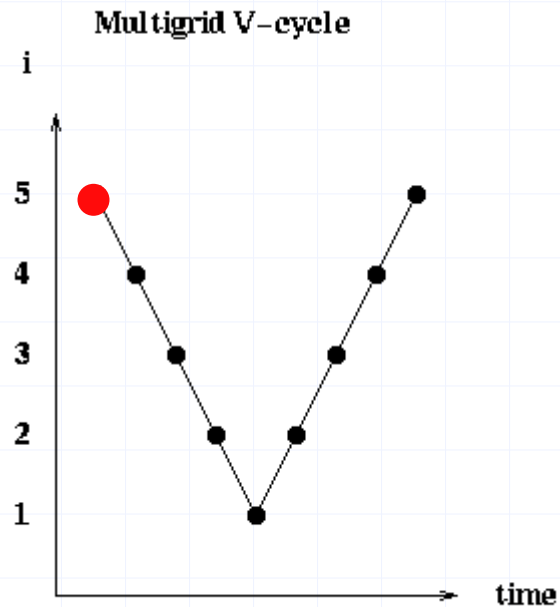
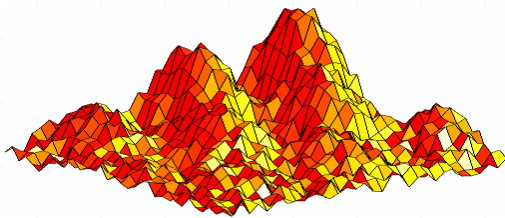
Must solve linear systems of the form

$$(\mathbf{K} - \omega^2 \mathbf{M}) \mathbf{x} = \mathbf{b}$$

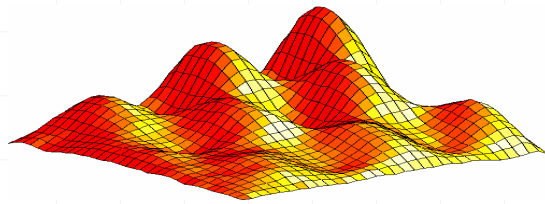


Geometric multigrid preconditioned GMRES
to non-Hermitian complex-valued linear system

Geometric multigrid

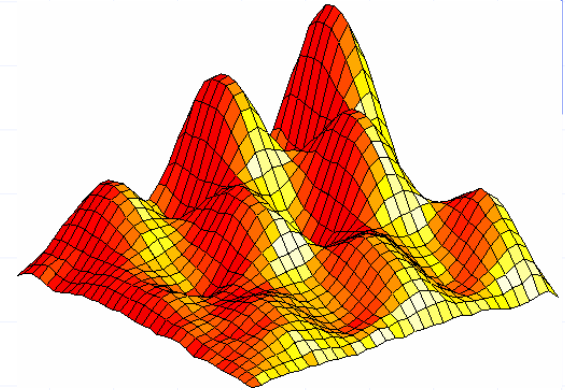
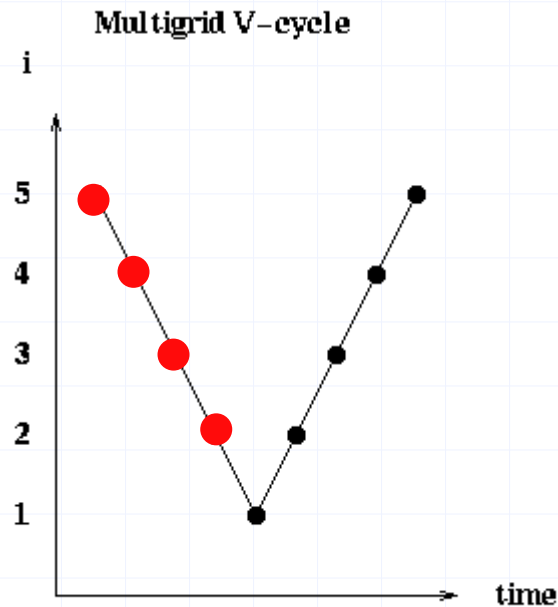


Geometric multigrid



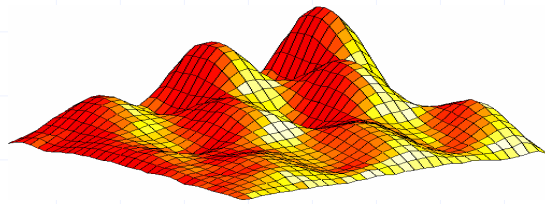
Smoother

Remove error with large eigenvalue



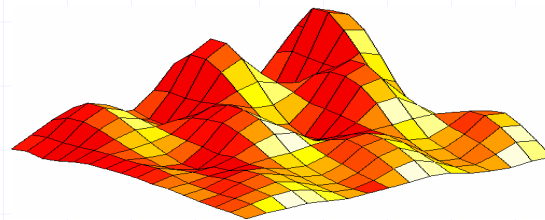
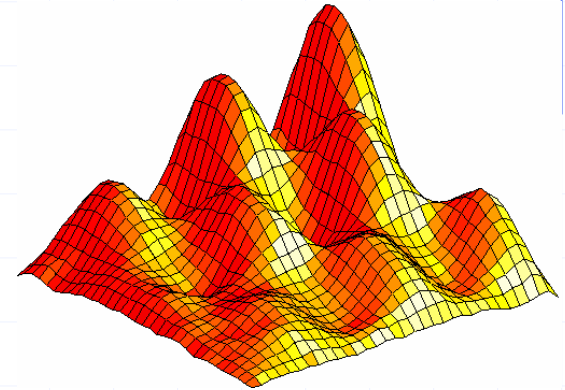
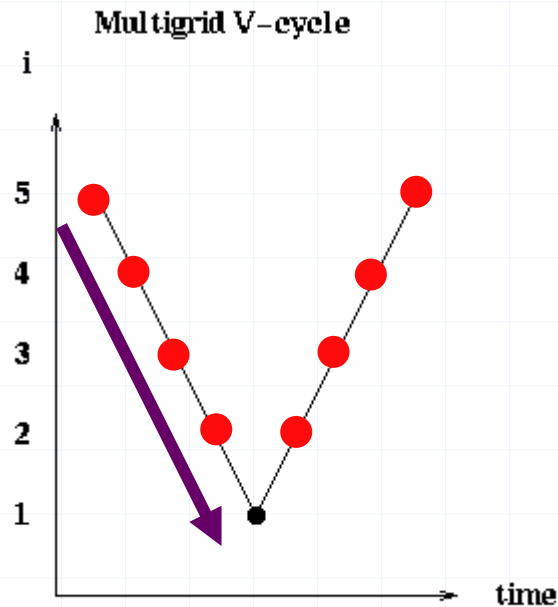
(Demmel, Parallel Computing Notes)

Geometric multigrid



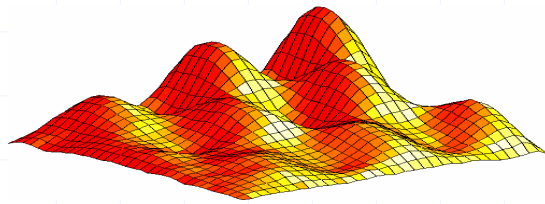
Smoother

Remove error with large eigenvalue



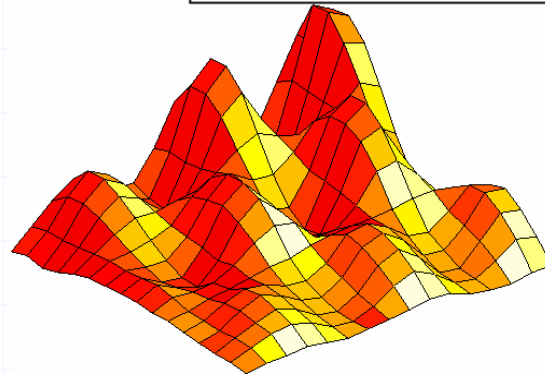
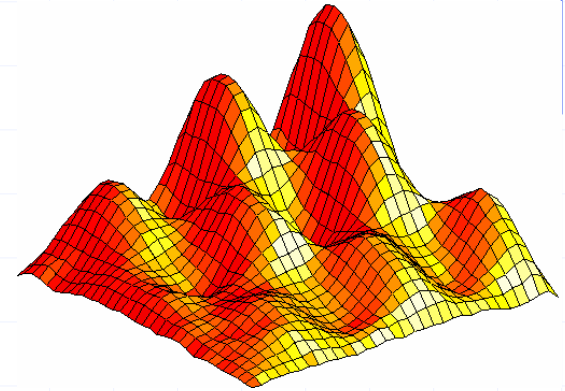
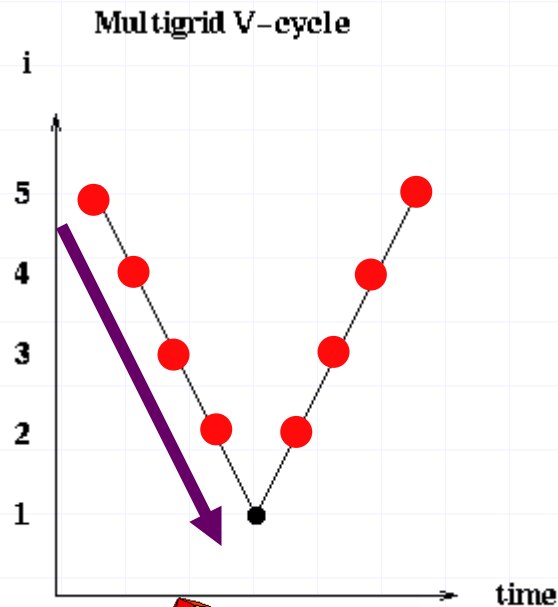
(Demmel, Parallel Computing Notes)

Geometric multigrid



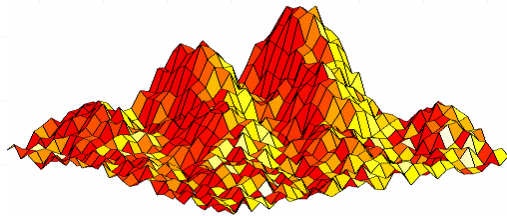
Smoother

Remove error with large eigenvalue



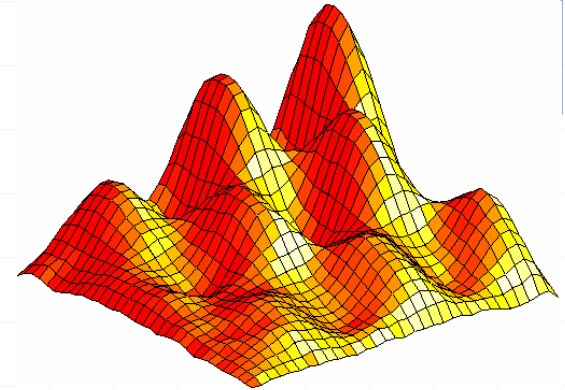
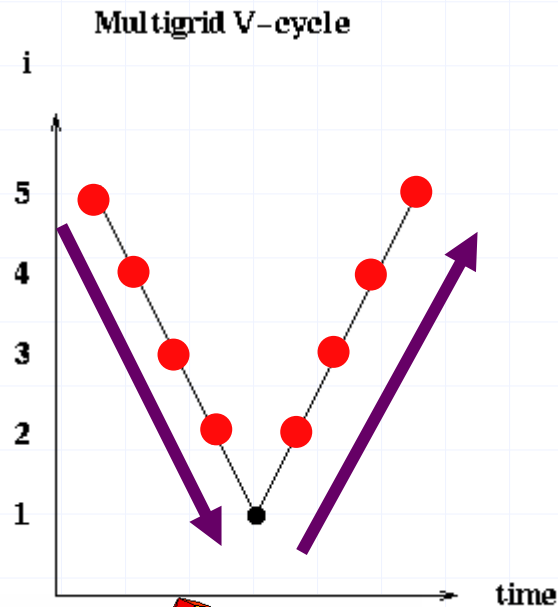
(Demmel, Parallel Computing Notes)

Geometric multigrid



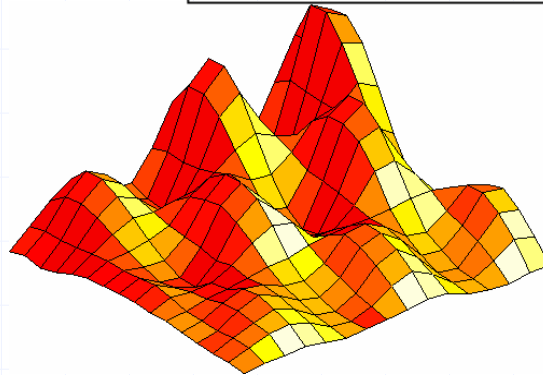
Smoother

Remove error with large eigenvalue



**Coarse grid correction
(Prolongation operator)**

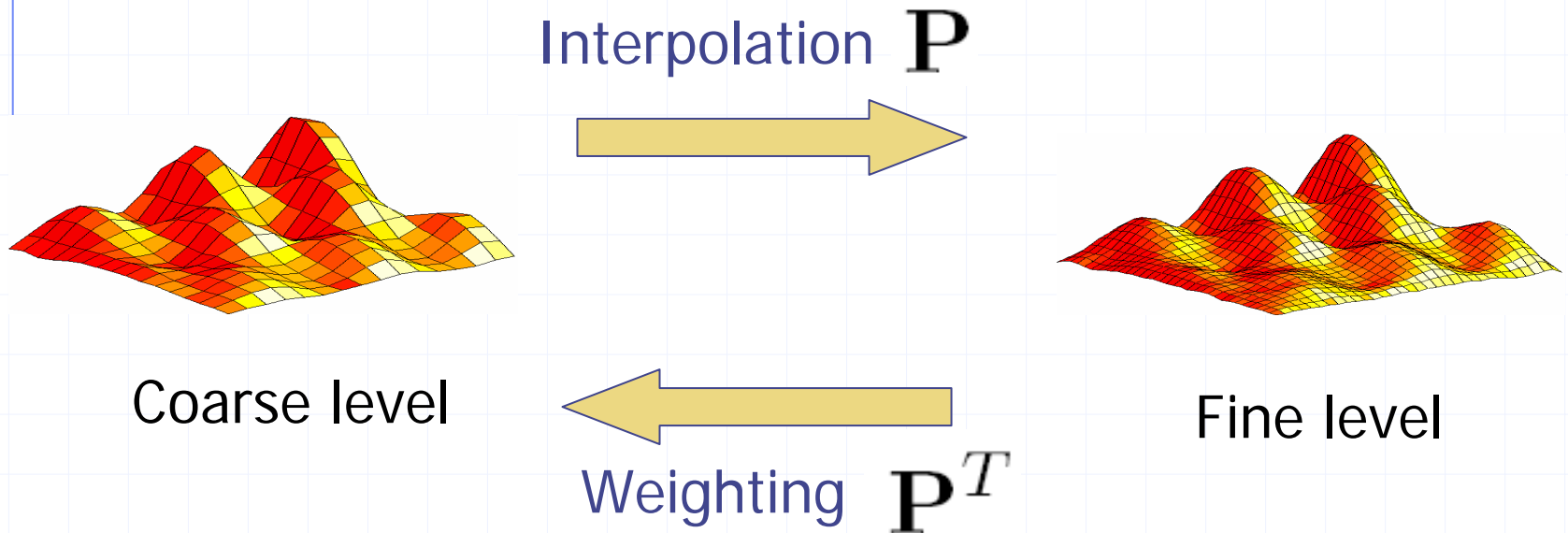
Remove error with small eigenvalue



(Demmel, Parallel Computing Notes)

Prolongation operators

- ◆ Maps solutions from one level to another with geometric information



Smoother for $\mathbf{Ax} = \mathbf{b}$

◆ No heuristics for choosing appropriate smoother for complex-valued non-Hermitian case

➔ Gauss-Seidel (Restriction in PML parameter selection.)

PML parameters

β : Absorption strength

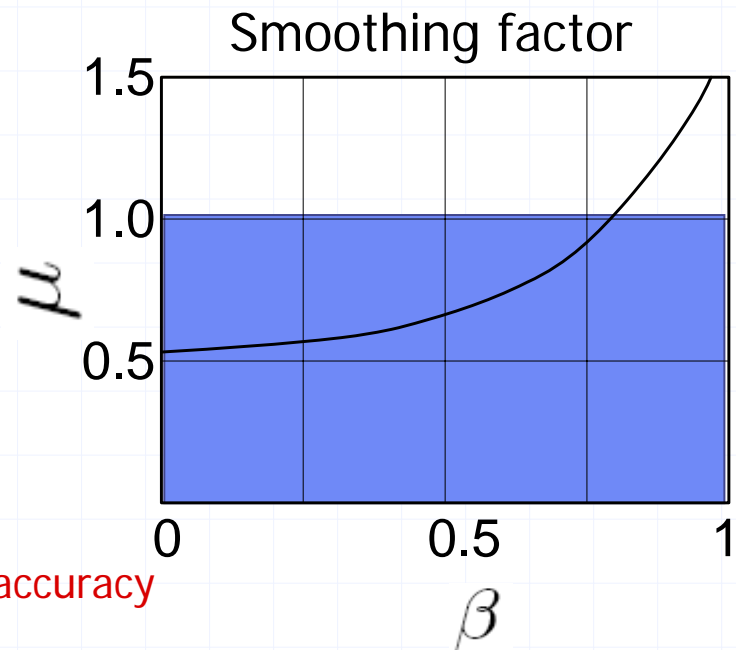
m : Wave lengths in PML

For convergent smoother, $0 \leq \mu < 1$

➔ β small Restriction for solvability

For small wave reflection

➔ m approx. 1-2 Restriction for accuracy



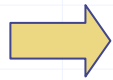
Two grid convergence factor

◆ Error reduction for multigrid with 2-levels

Fine level: n per wave

Coarse level: $n/2$ per wave

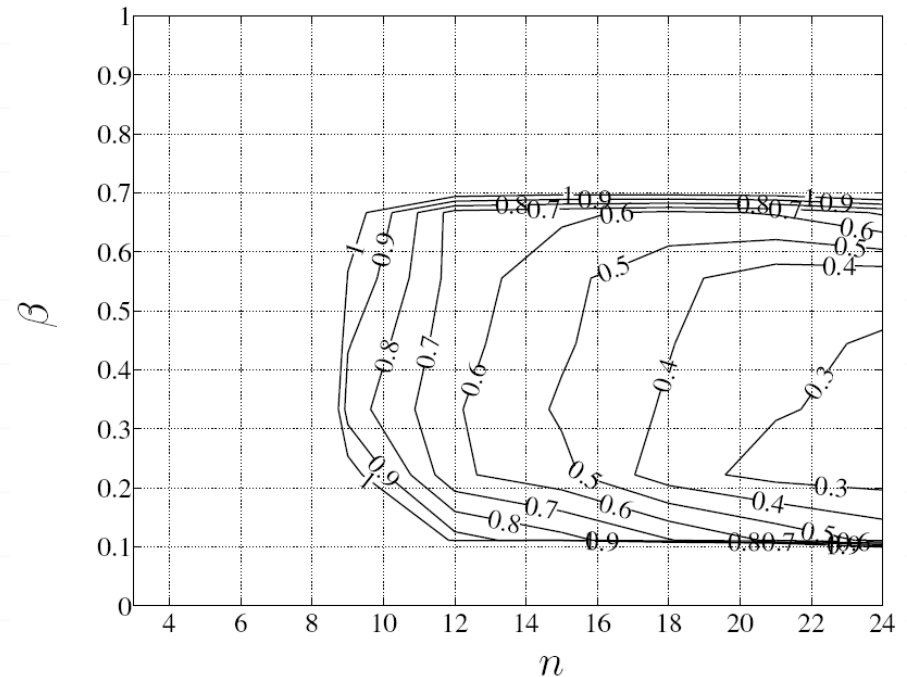
β : Absorption strength
 n : Nodes per wave on fine grid (discretization)



β at most 1
At least 6 nodes per wave on coarse grid

Restriction for multigrid convergence

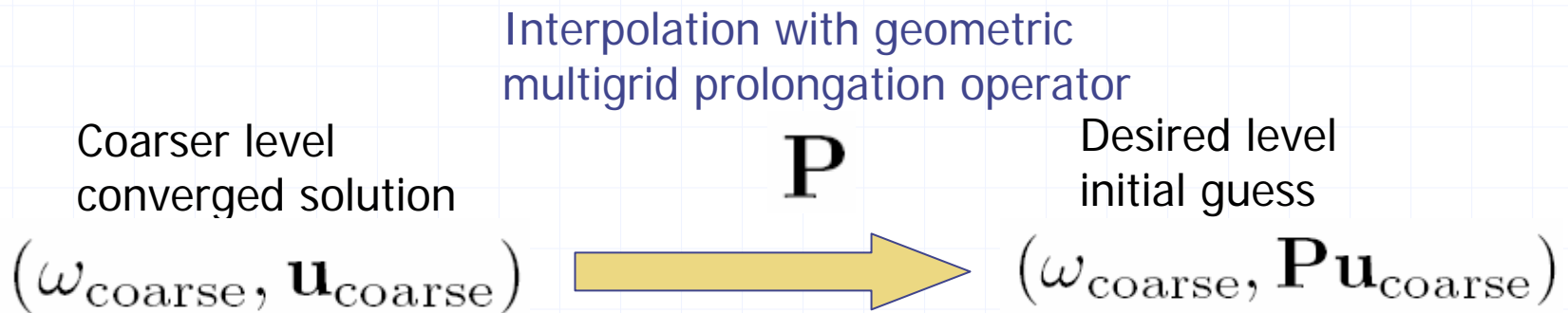
Gauss-Seidel



JDOQZ + GMG preconditioned GMRES

- ◆ The performance of JDOQZ depends strongly on initial eigenvalue/eigenvector approximation

➔ Use coarse grid solutions as initial starting points.

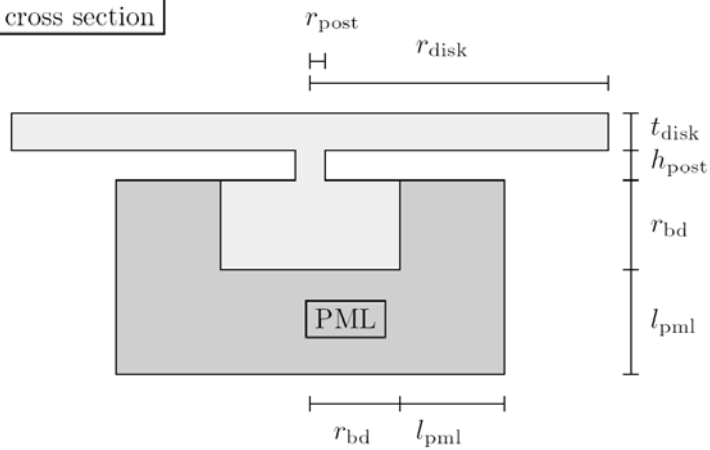


JDOQZ converges in 1-2 iterations (correction equation solves) to tolerance $1e-10$ with good initial guess. Scalability of iterative method is crucial.

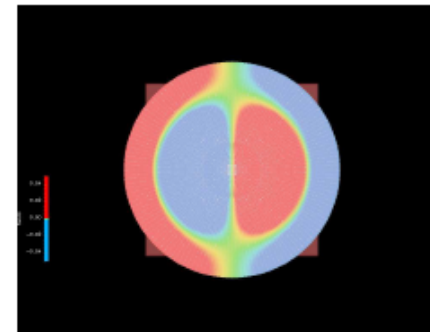
Numerical example: Disk resonator

◆ 2nd radial contour mode at 715 [MHz]

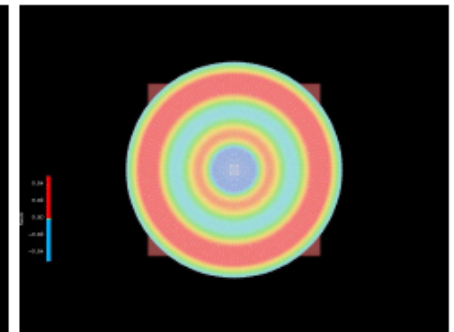
Resonator cross section



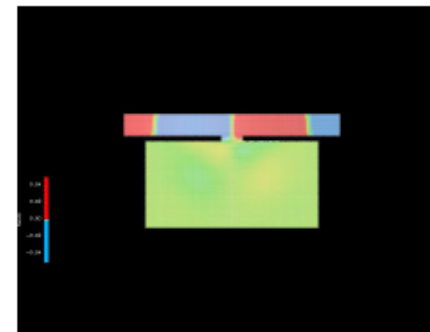
Top view: x displacement



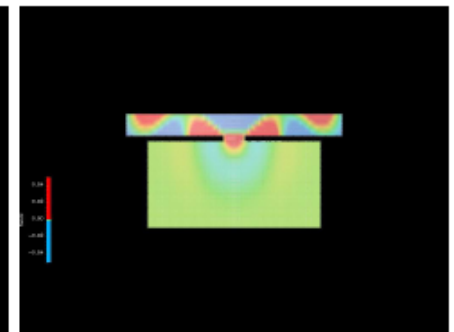
Top view: z displacement



Side view: x displacement



Side view: z displacement



Scalability of iterative method (PETSc implementation)

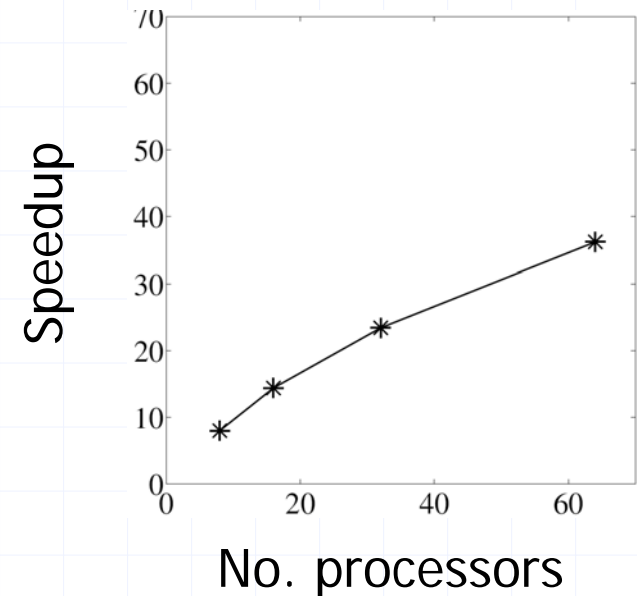
◆ Scaling with respect to size

◆ Speed-up with respect to number of processors

Level	Number of DOFS
1	49938
2	197574
3	977115
4	6140520

GMRES iters. for preconditioned residual of 1×10^{-10} for $\beta = 1.0$

Levels	1-2	1-3	1-4
715.7[MHz] (2nd mode)	49(1.2e-09)	51(1.6e-09)	41(3.3e-09)



Conclusions

- ◆ An algorithm for solving complex-symmetric eigenvalue problems arising from mechanical systems with Perfectly Matched Layers (PML)
- ◆ Jacobi-Davidson QZ + geometric multigrid
 - scalable, fast,
 - ◆ converging in 1-2 iterations with coarse grid initial values
 - restrictions on PML parameters
 - ◆ converging smoother: β small
 - ◆ converging multigrid : at least 6 nodes per wave on coarsest grid
 - ◆ numerical accuracy : approx. 1-2 wave lengths in PML